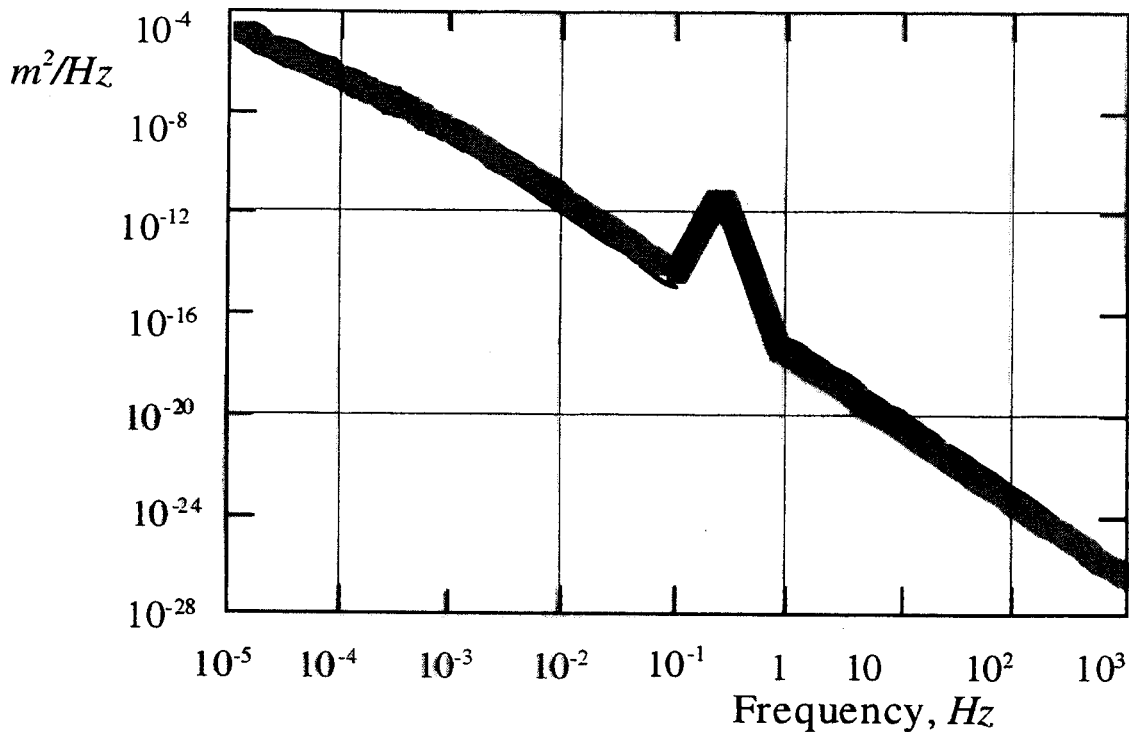


# **GROUND MOTION MODEL FOR THE LHC**

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# 1 INTRODUCTION

Power spectrum



(cultural noise taken out)

- power of ground vibrations increases steeply with decreasing frequency

====> possibility of non-negligible orbit deformations if the motion of accelerator quadrupoles is uncorrelated

- what about plane wave excitation ?

ground motion wavelength matches the betatron wavelength

====> very **narrow band** effect for frequencies around 1 Hz, *small spectral power*, small separation effect in LHC < 1/1000 of *rms* beam size

## 2.2 Basic model

The model consists of a transfer function and a source of excitation

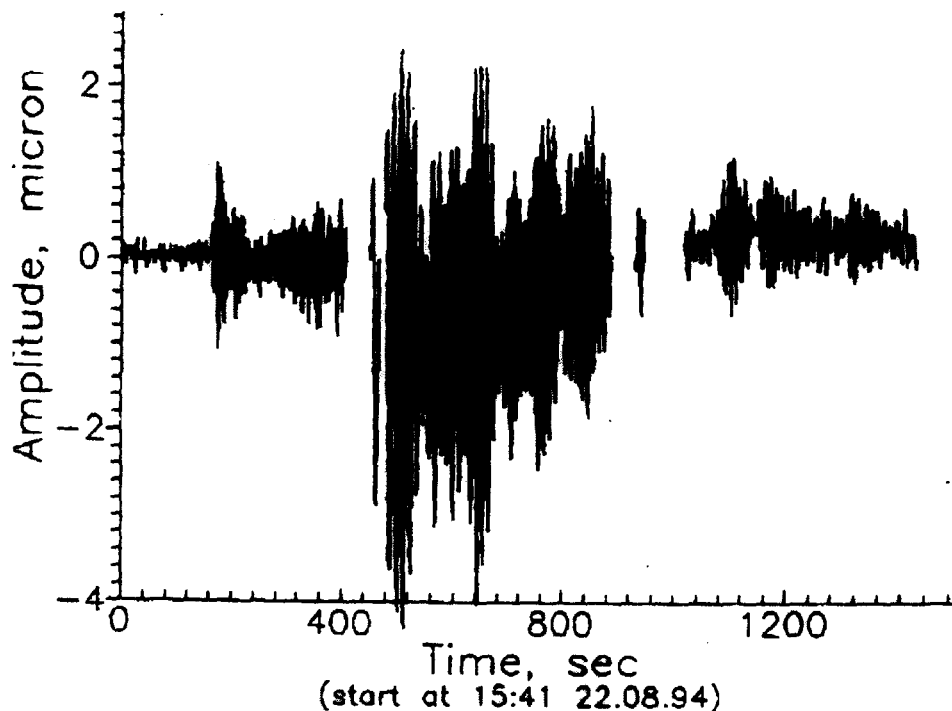
### 2.2.1 transfer function

maximum seismic length of the earth is  $\sim 1500$  s

seismic 'depth' of the earth is about 1/3 of this

$\implies$  a cut-off frequency of  $f_{co} \sim 2$  mHz

confirmed by the far away amplitude response of earth quakes:



fast oscillation is the response of the oceans pounding on the continents

amplitude response is not incompatible with a high-pass filter behaviour with a power cut-off frequency  $f_{co} \sim 1.5$  mHz

note : seismic wave attenuation with distance is very small

## 2.2.2 source

the obvious source is : *earth quakes!*

the transfer model suggests a source with a  $f^{-3}$  frequency slope to reproduce noise spectra

examination of a substantial body of phenomenological material has led to the *Gütemberg- Richter law* :

$$\log(n) = -M,$$

$n$  is the number of earth quakes in a given area with a magnitude  $M$  or larger

$$M = \frac{1}{m} \log\left(\frac{x}{x_0}\right)^2$$

$x_0$  is the lower observation limit of seismographs.

According to *Hamblin*  $m = 2$  so that :

$$\left(\frac{x}{x_0}\right)^2 = \frac{1}{n^2}.$$

The power density of the source can be found by differentiation :

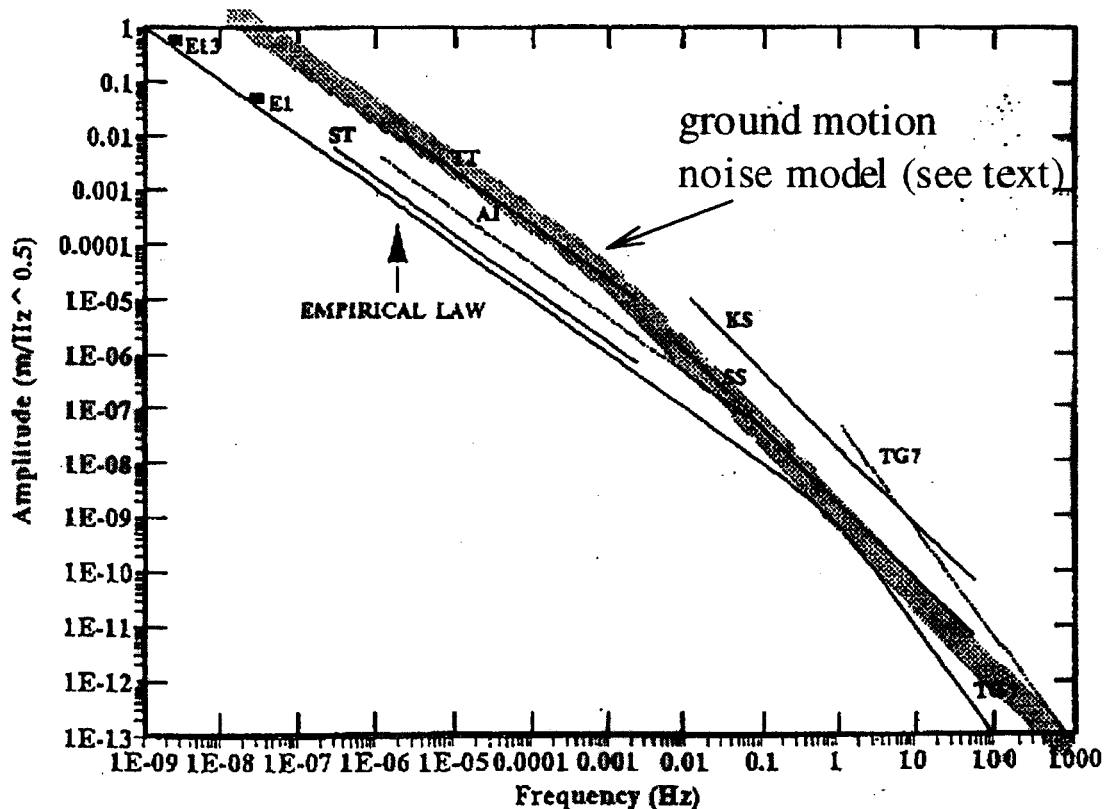
$$\frac{d(x/x_0)^2}{dn} = -\frac{2}{n^3} \propto -\frac{1}{f^3}$$

Combining transfer function and source of vibrations yields the ground motion spectrum :

$$\frac{dx^2}{df}(f) = \frac{k_{gm}}{f^2 \sqrt{f_{co}^2 + f^2}} [m^2/Hz].$$

$k_{gm}$  is a proportionality factor of the source to fit the observed spectral power density. It is a non-local quantity that varies from  $\sim 10^{-18} m^2/s^2$  to  $\sim 10^{-16} m^2/s^2$  depending on the state of global excitation

Comparison with various observations :



The line marked 'empirical law' is related to the model proposed by Takeda et al.

### 3 OBSERVATIONS WITH BEAM

coherence length can be determined from orbit measurements assuming an average value for  $k_{gm} = 10^{-18} \text{ m}^2/\text{s}^2$ .

integration noise spectrum yields displacement<sup>2</sup> of a single element :

$$dx^2(t) = \frac{k_{gm}}{f_{co}^2} \left( \sqrt{1 + (f_{co}t)^2} - 1 \right).$$

orbit deformation : multiply with an optical amplification factor :

$$O_A = (\beta Kl / 2 \sin(\pi q))^2 N$$

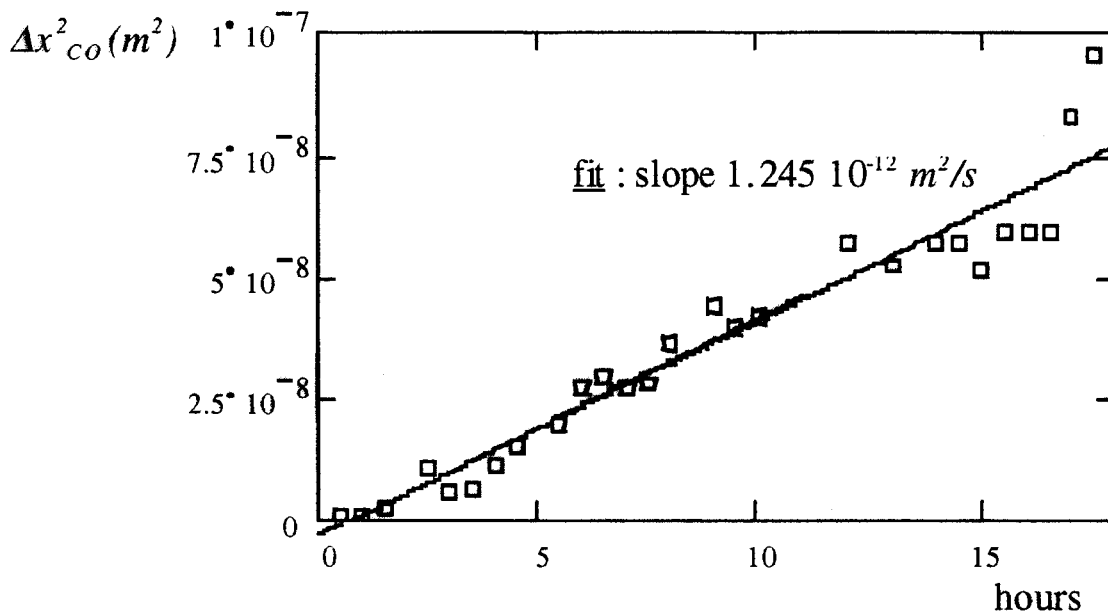
$\beta$  the optical function at a quadrupole  
 $Kl$  integrated focalisation force  
 $N$  is the number of **uncorrelated** blocks around the accelerator (at the maximum the number of  $F$  or  $D$  quadrupoles)

Brinkman & Rossbach (1994)

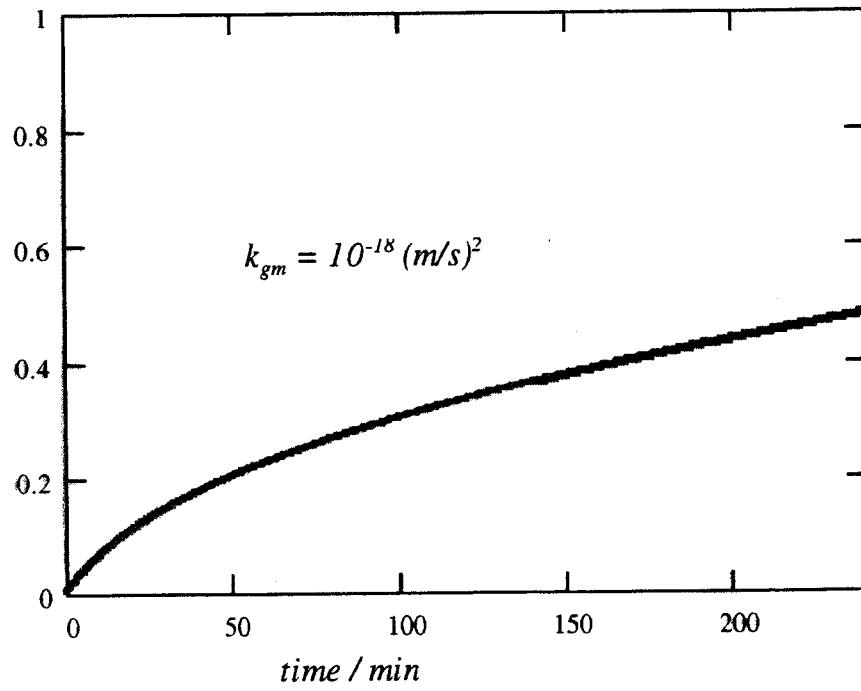
HERA-proton  $L_{ch} : 250 \text{ m}$   
 HERA-electron  $L_{ch} : 280 \text{ m}$

Tecker (1996)

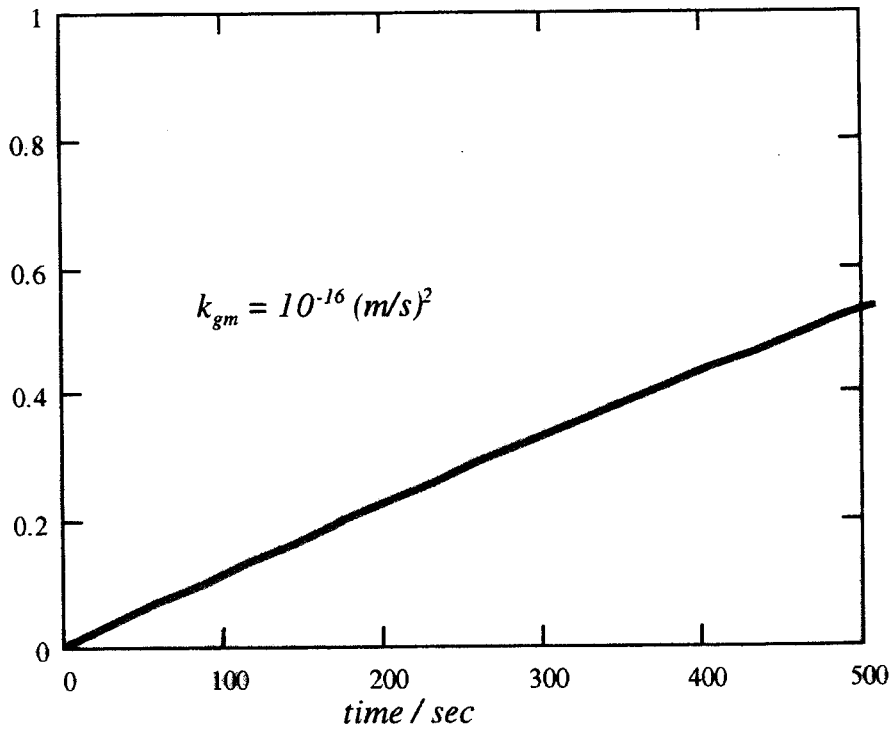
LEP  $L_{ch} : 130 \text{ m}$   
 (known effect supra quads removed)



half sep. /  $\sigma$



half sep. /  $\sigma$



Problem with the ATL interpretation :

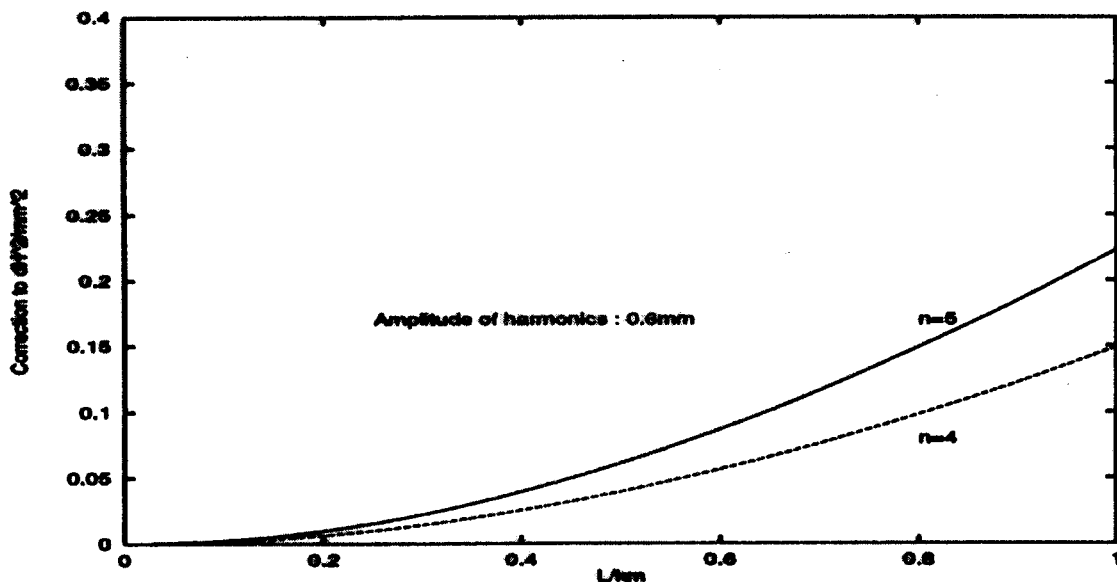
Function  $f(s)$  subtracted from the positions  $y(s)$

$$dH^2(L) = \int_0^C [y(s+L) - f(s+L) - y(s) + f(s)]^2 ds$$

- the functions  $f(s)$  and  $y(s)$  are uncorrelated
- the integral of  $y(s)$  is zero

$$dH^2(L) = dH_0^2(L) + \int_0^C [f(s+L) - f(s)]^2 ds$$

Harmonic :  $f(s) = a \times \sin(2\pi ns/C)$





# **TO BE KEPT IN MIND FOR THE LHC**

## **Short term :**

- **Normal level of excitation : “manual” control**
- **High level of excitation : orbit feedback**

## **Long term :**

- **Alignment data in LHC database for evaluation of annual realignment**
- **Systematic analysis of the closed orbits to detect important motions**