NLC -Note #24
April 1997

Erata to NLC-Note #24

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R.M. Jones (4- 14-97).

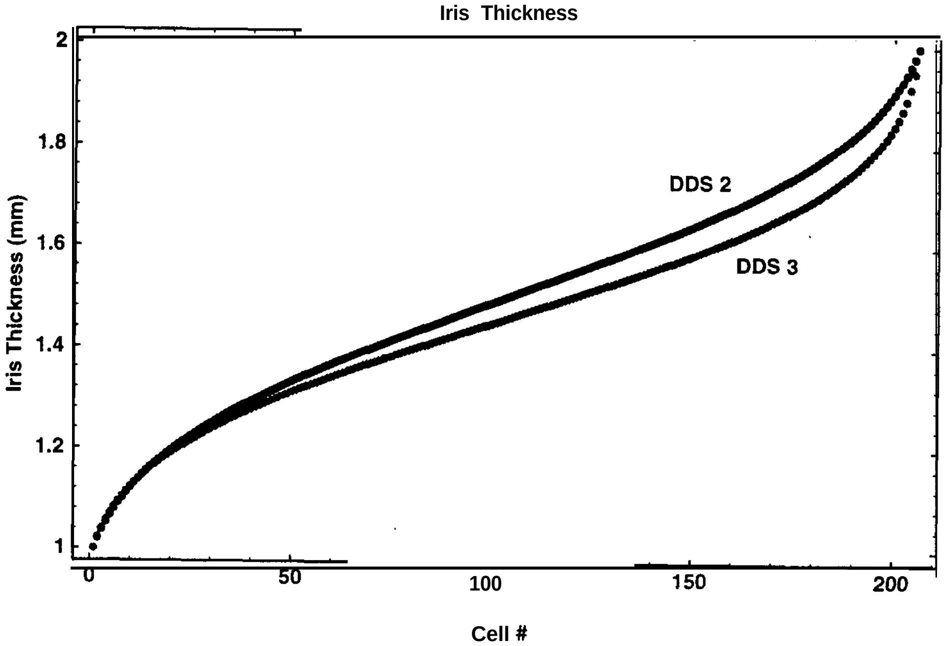

FIGURE 5 c: Cell geometrical parameters: iris thickness, t.

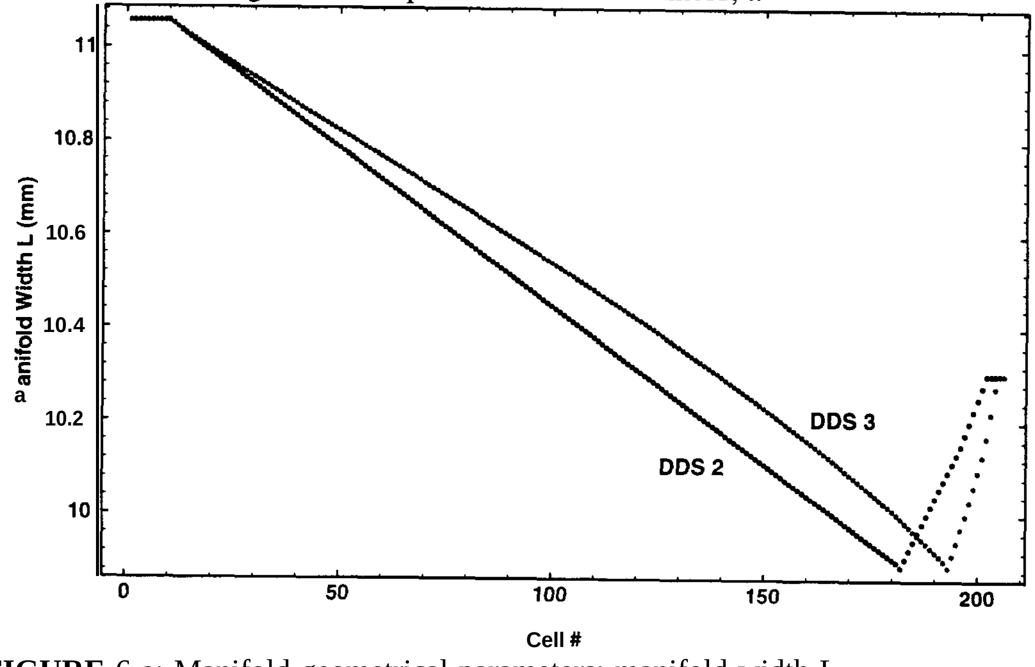

FIGURE 6 a: Manifold geometrical parameters: manifold width L.

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A Rapid Re-Design Method for the NLC DDS and Phase Stability of the Accelerating Mode: Implementation and Results

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Abstract:
In order to reduce the dipole wake encountered by the first few bunches accelerated in a multi-bunch NLC scenario the DDS (damped detuned structure) was re-designed such that a much improved Gaussian fall-off occurs in the initial wake-function. From the 9 parameterised model of DDS 1 we use a mapping function to allow DDS 3 & 4 to be modeled and hence avoid additional and prohibitively time consuming MAFIA runs. The equivalent circuit parameters and geometrical parameters are treated as functions of the synchronous frequency and are readily mapped onto the new synchronous frequencies. The new geometrical parameters form a family where each is associated with the iris diameter. In order to assess the deviation of the phase from $2\pi/3$ the difference between the new and old parameters is investigated. In many cases the difference is found to be of the order of a few tens of nanometers

1. INTRODUCTION

The first ever manifold DDS was designed such that the geometrical parameters (iris radius and cavity radius) of the cells were inverse functions of error functions, Erf. Further, the mode density function (the reciprocal of the derivative of the uncoupled frequencies with respect to mode number: dn/df) was prescribed to be Gaussian. However, the short range dipole wake function, is given by twice the inverse transform of the dipole kick-factor (K) weighted density function and under a Gaussian prescription of dn/df, 2Kdn/df is markedly asymmetric. The consequence of the asymmetry is a poor definition of the minima in the short range wake function. DDS 3 & 4 have been re-designed under a Gaussian 2Kdn/df prescription with a bandwidth of 4.71 units of sigma (with sigma 2.125% of the central frequency of the Gaussian) and this leads to a significantly improved short-range wake function.

The inverse Fourier transform of the spectral function [1] allows the global wake-function to be evaluated. However, in order obtain the new spectral function, all nine parameters of the structure must be obtained for 206 cells. This is a substantial computational task in running the MAFIA code required for the spectral function and in the careful fitting procedure required for all the new functions. However, the method used herein obviates this excessive computational work and requires that we fit all the 9 circuit model parameters together with the beam kick-factor to ten functions which all depend on the synchronous frequency only. The new set of new synchronous frequencies, dictated in our case by the requirement that $2Kdn/df$ be Gaussian, allows the 206 x 10 new characteristic parameters to be calculated. Similarly the 5 parameters, which define the geometry of the structure (the iris radius a, cavity radius b, iris thickness t, the radial distance of the edge of the manifold from the center of of a cell, H and the height of the manifold L) are also functionally dependent on the synchronous frequency of the beam and, under a new set of frequencies, 5 x 206 new dimensions are calculated for the DDS. Thus, in order to obtain the wake function and new geometrical parameters for fabrication all that is necessary is to obtain 15 functions.

However, under this new mapping one might express some concern as to whether the properties of the fundamental (i.e. accelerating) mode have been adversely affected and so with this in mind we conducted an intensive investigation as to the deviation of the cell dimensions, parameterised by the cavity diameter 2b, from their values designed in DDS 2 (all dimensions form an invariant family parameterised by the cavity diameter b). This is detailed in section 4 and successive sections.

2. THE MAPING FUNCTION

In our design of DDS 2 we chose eleven representative sections to obtain frequency-phase pairs from detailed MAFIA simulations and hence obtain ten model parameters (nine circuit parameters plus the cell kick-factors) for each of the eleven sections. Parameters for all sections are subsequently obtained by error function fits and interpolation. A similar procedure may be followed to determine the five geometric parameters (i.e., cell and manifold dimensions) for all the sections from those for the original eleven. This is a substantial task for each structure design. However, as we now have all fifteen parameters as a function of synchronous frequency, we can take advantage of this functional dependence to explore new design distributions and to obtain the set of section dimensions which would be needed to realize them.

Based on our fit parameters we prescribe a smooth uncoupled spectral function, $S_0(f_s)\lambda$ and impose the condition that: $2K(f_s)dn / df = S_0(f_s)\lambda$, where K is the uncoupled kick factor, f_s the synchronous frequency and, λ is a scale factor to be determined. The upper and lower truncation bounds on the synchronous frequencies are imposed, f_{s1} and f_{sN} and the normalisation condition is obtained:

$$\lambda = N \Big/ \int_{f_{s1}}^{f_{sN}} (\tfrac{1}{2} S_0 / K) df_s \tag{2.1}$$

Then the new synchronous frequencies are determined according to:

$$\int_{f_{sn}}^{f_{sn+1}} \frac{1}{2K} \lambda S_0 df_s = 1 \tag{2.2}$$

This enables all the cell synchronous frequencies to be determined and hence the new ten parameters are determined.

This procedure is implemented in the following section to calculate the spectral function and associated wake function for DDS 3 and 4. In the proposed design DDS 3 and 4 will be formally identical and so hereafter when we refer to DDS 3, DDS 4 is also included by implication.

3. CALCULATION OF THE WAKE FUNCTION

In the revised design for DDS 3 we chose a truncated Gaussian distribution for the uncoupled $2Kdn/df_{syn}$ distribution, with a bandwidth of 4.71 units of sigma, (a bandwidth of 10.159% of the central frequency and sigma is 2.125% of the central frequency) and this provides a basis for the determination of the 206 synchronous frequencies. The kick factor weighted density function for DDS 3 and DDS 1 are shown in Fig 1.

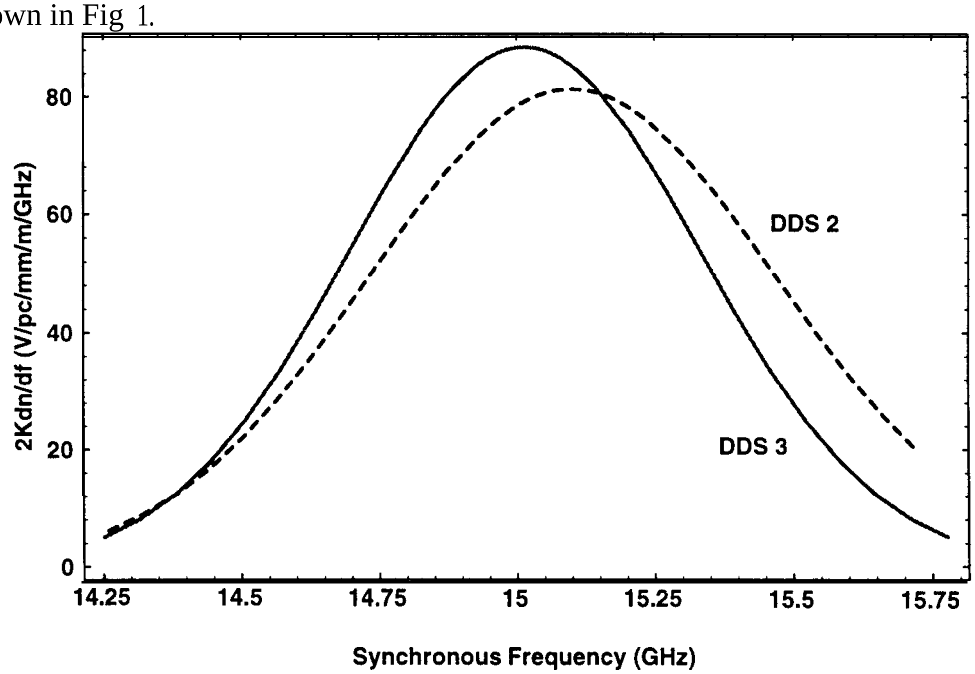

FIGURE 1: Twice the kick factor weighted density function for DDS 2 (shown dashed) and the corresponding function for the re-designed DDS 3 and 4

It is evident that DDS 2 is markedly asymmetric and this adversely affects the sharpness and depth of the minima for the short range wake function as will be seen in the subsequent calculation.

In order to calculate the wake function we first are required to calculate the spectral function associated with the 9 mapped parameters. This spectral function calculated for DDS 3, and shown in Fig 2, maintains the Gaussian characteristics imposed upon it from the synchronous frequency distribution but modulated with oscillations of large amplitude (resulting in a large part from reflections occurring in the higher order mode couplers in the manifold).

Also, the spectral function exhibits the underlying damped mode structure as mentioned previously for DDS 2 [1] and is shifted with respect to the 2Kdn/df curves, as given in Fig. 1. The difference between the respective curves in Fig. 1 and Fig. 2 becomes more pronounced for higher frequencies and this is a consequence of the synchronous frequencies becoming more highly perturbed as one progresses down towards the higher energy end of the structure where the coupling to the manifold has been designed to be largest. These coupled synchronous frequencies will be discussed in a future publication [2]

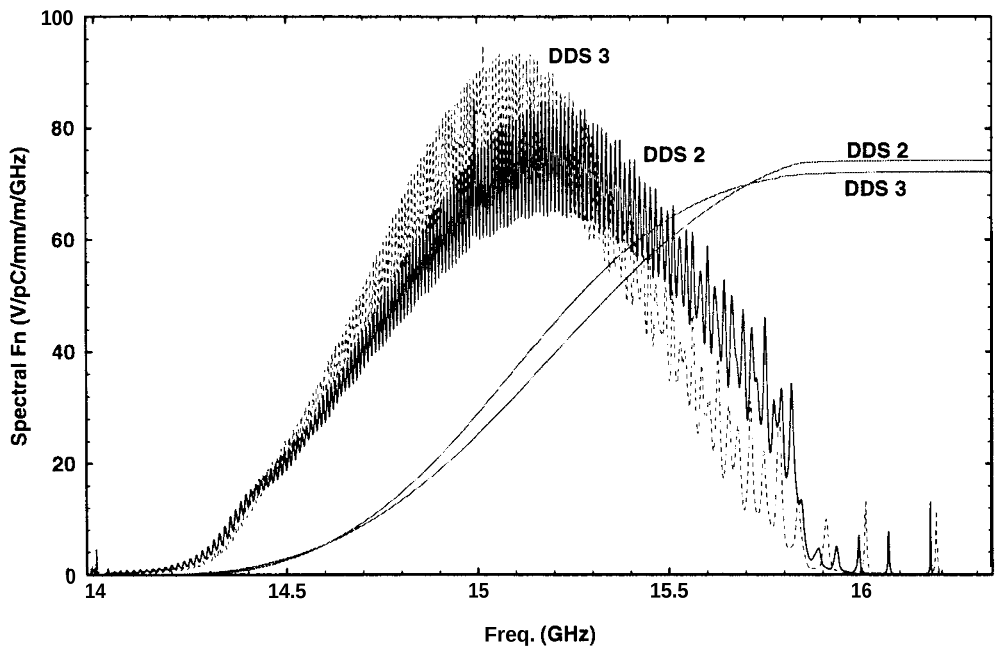

FIGURE 2: Spectral function for DDS 2 and DDS 3

It is important to note that the sharp, rather precipitous, fall-off in the spectral function in DDS 2 at approximately 15.8 GHz has detrimental affects on the short range wake function. In DDS 3 the spectral function falls off smoothly and gradually and this has beneficial effects on the range wake function in that it enables a faster fall-off to occur. Indeed for a perfectly smooth termination, which we refer to as our idealised case [1], it is possible to achieve more than an order of magnitude weaker wake function at the 90 bunch point.

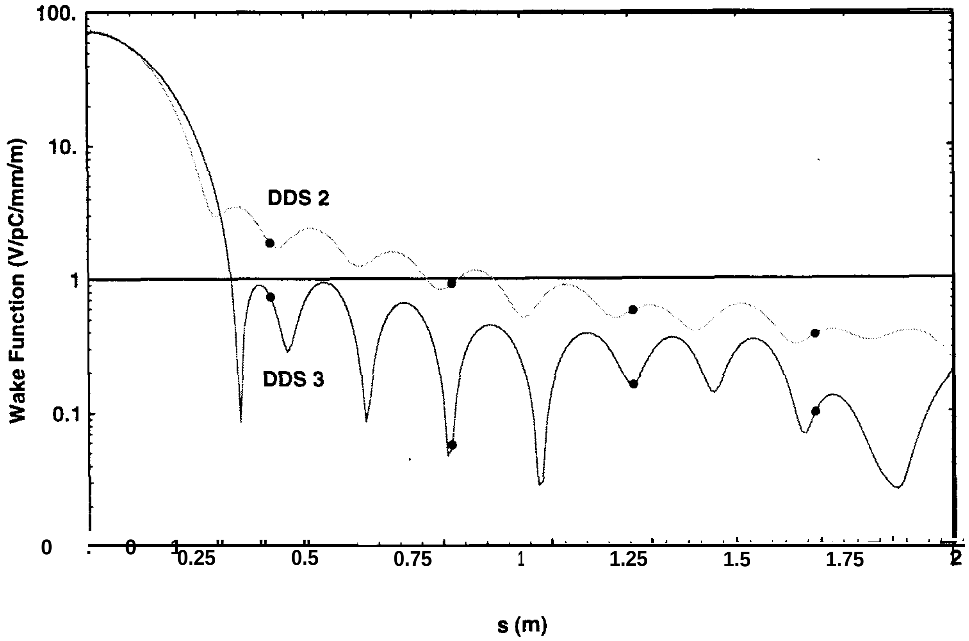

FIGURE 3: Short range wake function for DDS 2 and DDS 3

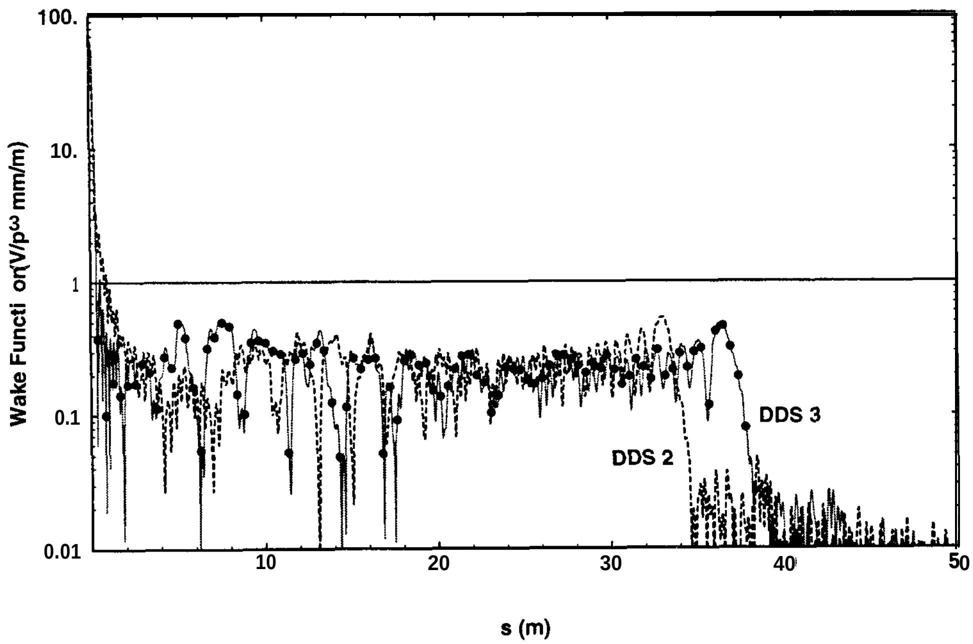

FIGURE 4: Long range wake function for DDS 2 (shown dashed) and DDS 3. The points are at the location of each of the bunches, of which there are ninety.

It is interesting to note that the maxima of DDS 3 is a little larger than that of its counterpart DDS 2. However, the area under the curve (bounded by the upper & lower synchronous frequencies for each structure) corresponding to DDS 3 is slightly smaller than DDS 2. This reduction in the area of the curves is indicative of a reduced wake function at the origin (since the wake function is given by the inverse transform of the spectral function) and this is in itself a consequence of the larger iris dimension in DDS 3.

The remaining parameters, viz, the those associated with the geometry of the DDS are outlined in the following section.

4. GEOMETRICAL PARAMETERS

Each of the five geometrical parameters are fitted with an interpolation function, the independent variable in each case being the synchronous frequency. Thence, armed with these new synchronous frequencies the new 206 x 5 parameters are readily obtained. Both the new and mapped cell parameters are shown in Fig. 5 and the manifold parameters in Fig. 6. The end points of DDS 2 and DDS 3 are identical by design so that no extrapolation is required in the determination of the DDS 3 parameters. Thus, in determining all parameters only third order interpolation between cell points has been employed, with a view to minimising any error in the generation of the new points.

It is evident from the curves that in the downstream end (or low energy end) of the structure the parameters are very close to that of the DDS 2 design whereas in the upstream end both the iris and the cavity diameter are increased significantly. This is a consequence of the asymmetry in the original design in which the kick factor weighted distribution reached too low a level in the upstream location of the structure and this has been corrected for in DDS 3. The thickness of the irises however, is reduced with respect to DDS 2, but this reduction is sufficiently small that the structure still maintains its mechanical integrity.

The manifold curves, shown Figure 6, illustrate the degree of tapering required to increase the TE coupling to the cells as one progress down the structure. This increased coupling is necessary because the modal composition of TE/TM is reduced as one moves towards the upstream end of the structure and to achieve a Q value in the neighborhood of a 1000 or so, increased coupling is required. There is a reverse in the taper towards the end cells in the upstream end and this is instituted in order to lower the cut-off frequency of the HOM (higher order mode) coupler and hence improve the match of the mitered bend of the HOMs at the lower frequency end of the band.

It is necessary to have well-matched HOM loads because the wake function is very sensitive to the power reflected back into the accelerator [1]. However, under the mapping the middle cell has effectively shifted forward by 11 cells and hence the upward taper (which occurs in cells 182 to 202 for DDS 2) maps the up-taper into a region where there are no manifold cells. Thus, in DDS 3 we change the

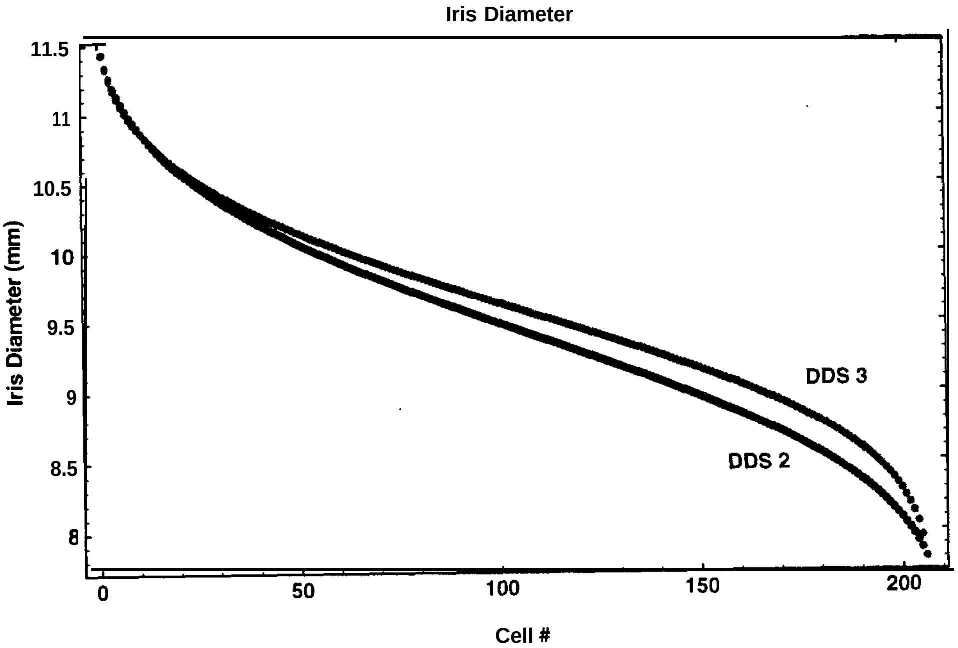

FIGURE 5 a: Cell geometrical parameters: iris diameter,

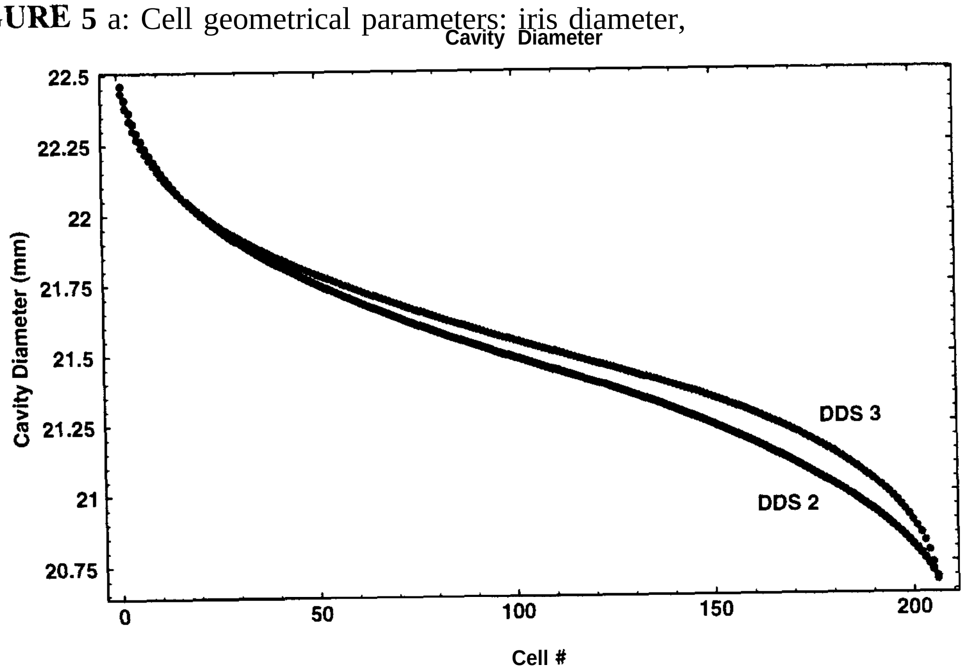

FIGURE 5 b: Cell geometrical parameters: cavity diameter

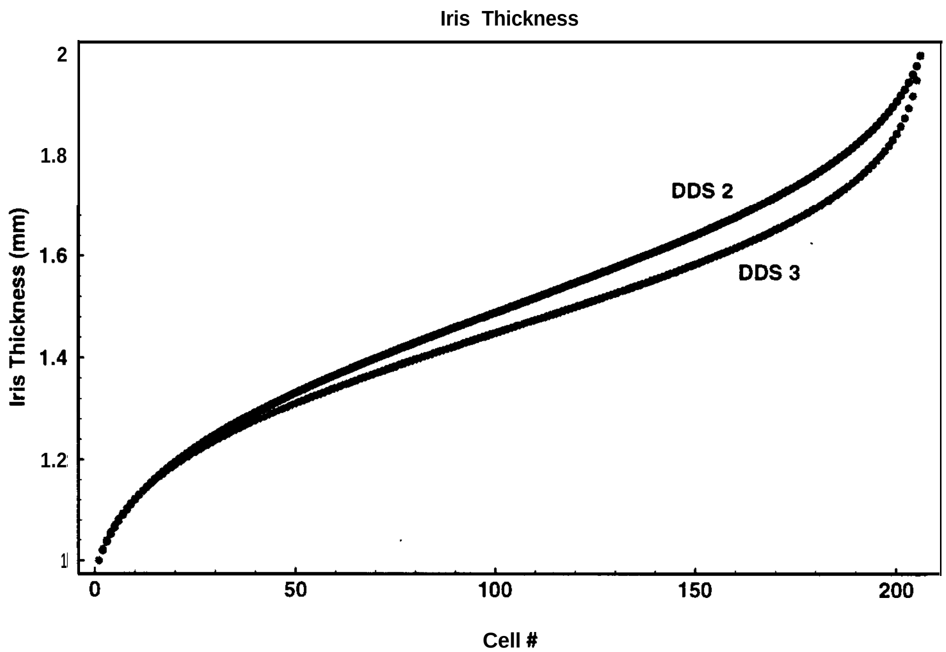

FIGURE 5 c: Cell geometrical parameters: iris thickness, t.

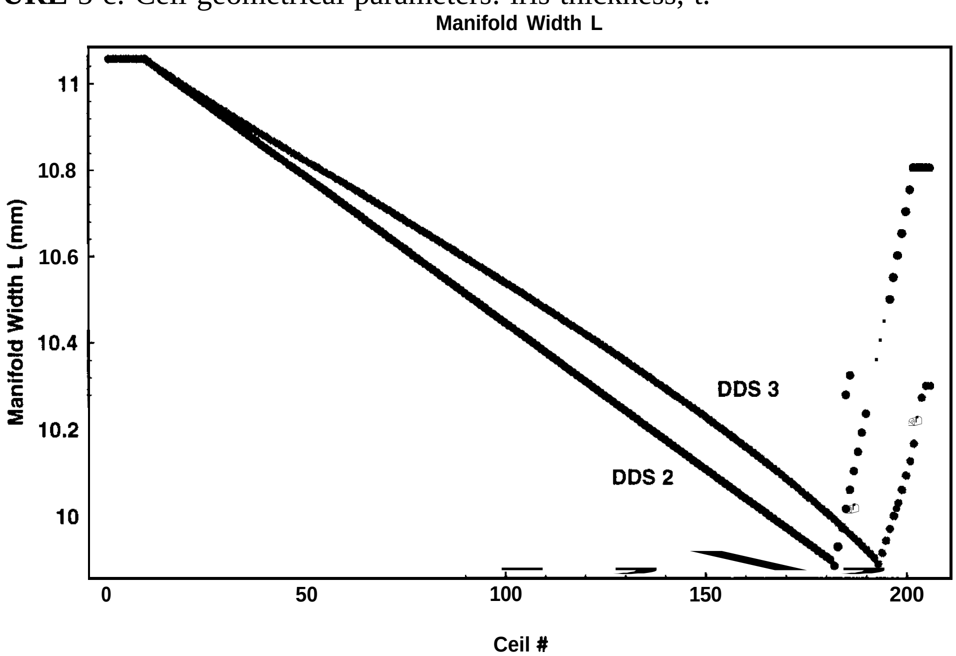

FIGURE 6 a: Manifold geometrical parameters: manifold width L.

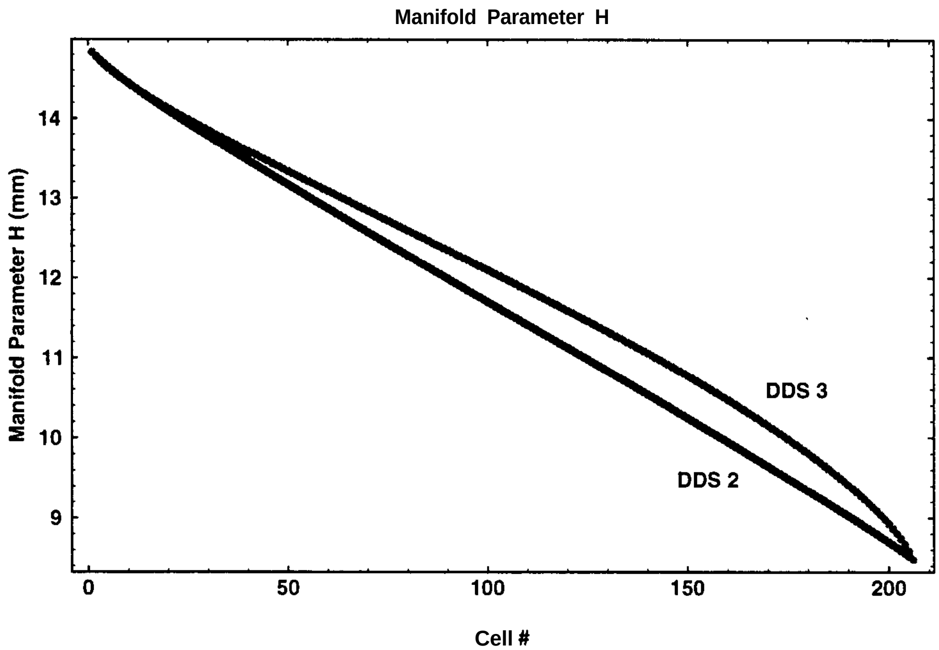

FIGURE 6 **b:** Manifold geometrical parameters: radial distance of the edge of the manifold the from center of a cell, H.

mapped taper in this region by increasing the gradient of the taper and withdrawing its to cell 202. This will adversely affect the wake function but we are confident that its deleterious effect will minimal.

However, it is important that the prescribed $2\pi/3$ phase of the accelerating mode be preserved and in this section a parameterised form related to phase dependence is investigated.

5. PHASE ADVANCE

All the geometrical parameters of the structure form a family functionally dependent on the iris radius a. This family is invariant under the mapping and hence the deviation of the parameters from those of DDS 2 serves as an indicator of how well the phase of the fundamental mode has been preserved under the mapping. The cell and manifold parameters are shown in Figure 7 and 8, respectively, with those of DDS represented by a continuous line and those of DDS 3 by dots. The iris thickness, t, for DDS 3 shown in figure 7a as a function of the iris diameter 2a, together with the difference between DDS 2 & 3 (inset) reveals the deviation to be no more than 5 nanometers. This is, of course more than adequate for our practical accelerator design. However, we also are required to calculate the deviation in the cavity diameter 2b and as unfortunately the

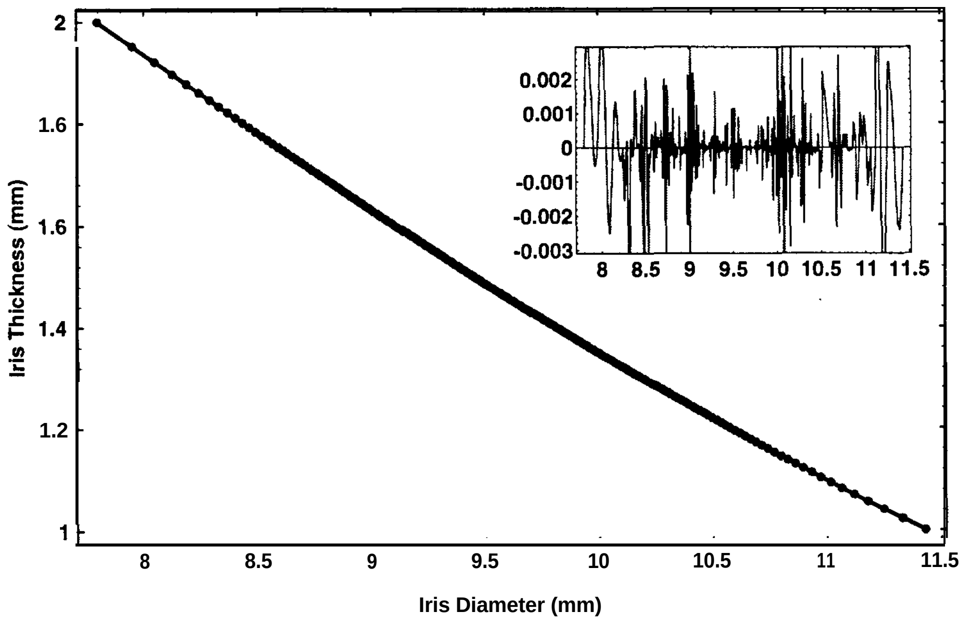

FIGURE 7 a: Iris thickness, t, as a function of the iris diameter, 2a. Also, shown inset, is the difference in μm between the original parameter t (corresponding to DDS 2) and that of the mapped parameter t (corresponding to DDS 3).

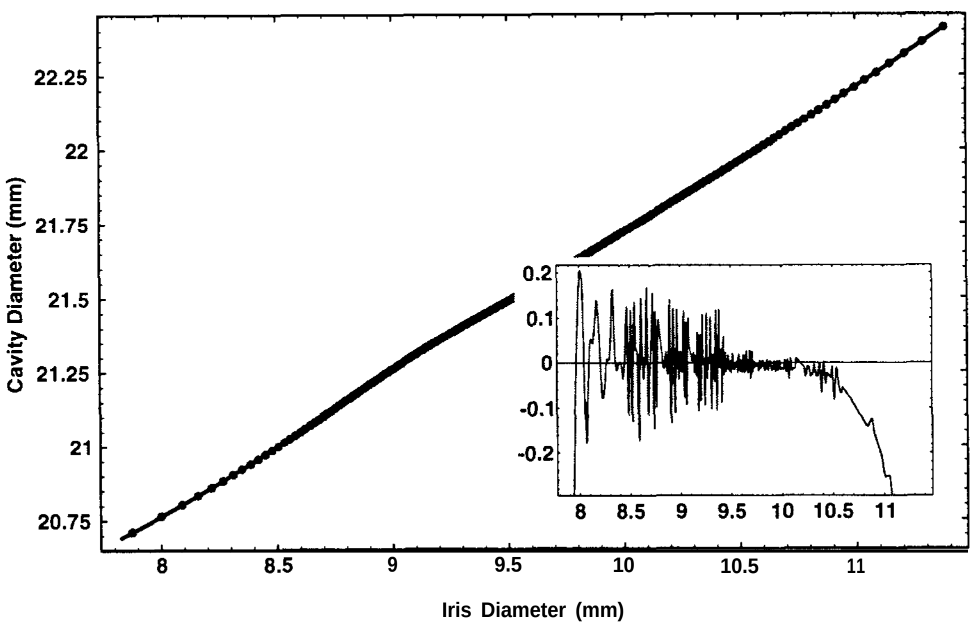

FIGURE 7 **b:** Cavity diameter, 2b, as a function of the average iris diameter, 2a. Also, shown inset, is the difference in μm between the original parameter 2b (corresponding to DDS 2) and that of the mapped parameter 2b (corresponding to DDS 3).

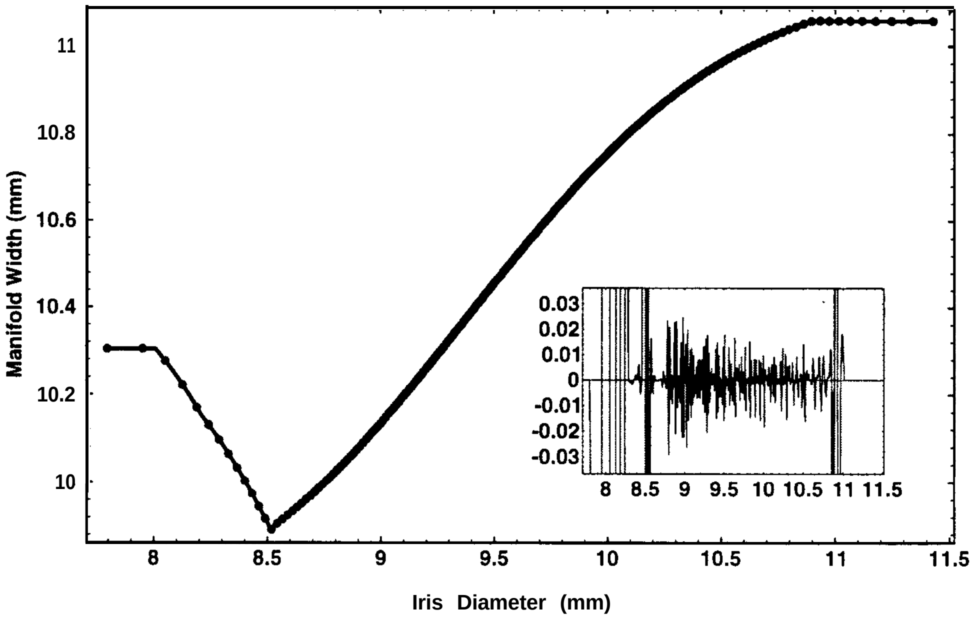

FIGURE 8 a: Width of manifold, L, as a function of the iris diameter, 2a. Also, shown inset, is the difference in μm between the original parameter L (corresponding to DDS 2) and that of the mapped parameter L (corresponding to DDS 3).

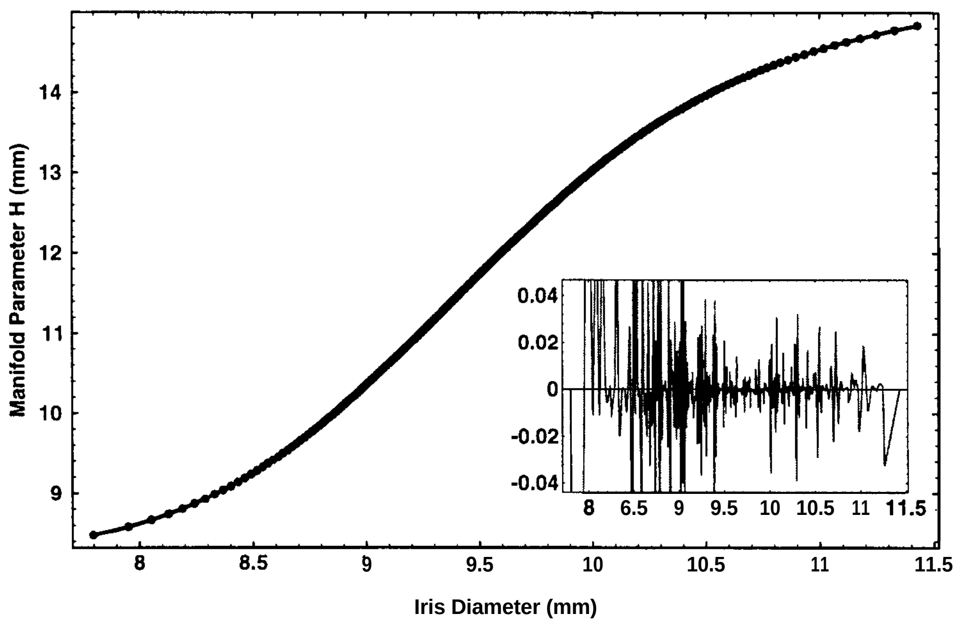

FIGURE 8 b: Distance of manifold from the center of a cell , H, as a function of the iris diameter, 2a. Also, shown inset, is the difference in μm between the original parameter H (corresponding to DDS 2) and that of the mapped parameter H (corresponding to DDS 3).

original design parameter are not readily available (only the values used in DDS 2 resulting from an averaging of the cups which compose adjacent cavities being available) we are required to reproduce the cavity diameter values prior to averaging. Averaging was removed by forming an interpolating function for the cavity diameters as a function of cell number and then selecting the half points as those of the un-averaged values. This process requires that the final point be extrapolated from the cavity data set and in the comparison of the data sets this region will be the most unreliable source of data. Nonetheless, performing this process and fitting an interpolation function through both the initial un-averaged cavity diameter data set, DDS 2 as a function of the averaged values of the iris radius (given directly by the data set used for fabrication) and, through DDS 3 as a function of the geometric average of the iris radius we obtain Fig. 7b. The deviation, shown inset to Fig 7b, over most of the range is approximately 0.15 μm and reaches 0.8 μm in the upstream part of the structure and this is considered to be acceptable for the design of our X-band accelerator. Similarly, the manifold parameter, L and H are remarkably stable with respect to the mapping as is illustrated in Figure 8, where the deviation is at worst approximately 50 nanometers.

6. CONCLUSIONS

We have developed a method to rapidly design new DDSs based upon a mapping procedure. This method enables both the wake function and the new geometry of the structure to be evaluated. Indeed, we have applied this method to calculate the short range wake function for DDS 3 and we find that, on average, the wake function is reduced by a factor of 5 or more. The new geometrical parameters form an invariant family which are functionally dependent on the cavity iris diameter, 2a and the deviation of the new family of parameters from the old provides an indication as to the accelerating mode's phase advance of the new structure. We discovered, in many cases, the discrepancy between the new and old parameters to be no more than a few tens of nanometers and at worst to be 0.8 μm in the end cells of the structure. This deviation is deemed to be sufficiently small that as far as is practically possible, imperceptible adverse affects are expected to be observed on the phase of the fundamental accelerating mode of the structure.

Furthermore, we have also used this technique to analyse the wake function in a structure consisting of groups of identical cells and we find that for groupings of 4 cells (reducing the total number of differing cells to 70 or so) that the wake function is seriously affected (to the extent that it is enhanced by a factor of 4 or more [3]).

7. ACKNOWLEDGMENTS

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